

Wave Optics

why focusing ?

- ⊕ sensitivity !
- ⊕ angular resolution and /or precise source identification
- ⊕ smaller detectors give better spectral resolutions
- ⊕ lower mass detectors, cooling systems, shielding
- ⊕ lower data transmission and storage
(data rate used for signal rather than for noise)

but how ?

reflection

refraction

diffraction

high energy astrophysics : $1/\lambda$ irony

The performance of virtually all astronomical instruments is diffraction limited

angular resolution	A	~	d/λ
spectral resolution	R	~	N/λ

In optical astronomy : $d/\lambda \approx 10^6\text{--}10^7$. Best angular resolutions are presently achieved in radio astronomy ... because wave properties are actually *used* rather than just presenting a limitation. VLBI over baselines (d) of up to $10^8\text{--}10^9$ wavelengths $\Rightarrow \Delta\theta \approx$ milliarcsec resolution

In high energy astronomy, even a system with an aperture of a millimeter should produce diffraction limited angular resolutions of the order of $> 10^9$ (1 MeV), yet present instruments have allowed for

angular resolution (gamma)	$\Delta\theta$	~	$0.1^\circ - 5^\circ$
spectral resolution (gamma)	R	\leq	500

focusing instruments : from *total external reflection* to *Laue diffraction*

total external reflection

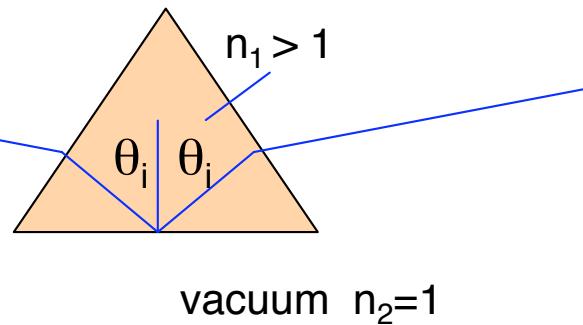
Wolter telescopes	tradl, replicated, foil optics	0.1- 12 keV
Lobster eye telescopes, Microchannel Plate Optics		0.1-3.0 (+) keV
Capillary Concentrators “polycapillaries”		1 - 60 keV
Multiple reflection optics (4+)		1- 80 keV
Kirkpatrick/Baez Optics		

diffraction

Multilayer mirrors	Unifrom Period Multilayers	
	EBB Multilayers	20 - 100(+) keV
Bragg-Lenses		10 - 200 keV
Laue-Lenses	broadband Laue lens	
	narrowband-, tunable- Laue lens	100 keV - 2 MeV
Fresnel zone plate		
Phase Fresnel lens		< 1 keV - > 1 MeV

total external reflection

total internal
reflection
of visible light

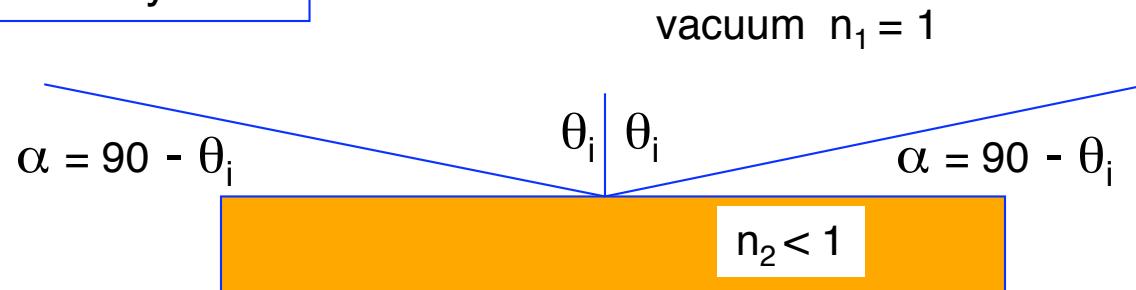


Snell's law

$$\sin \vartheta_2 = \frac{n_1}{n_2} \sin \vartheta_1$$

no solution for
 $n_2 < n_1 \sin \theta_i$

total external
reflection
of X-ray



Snell's law

$$\sin \vartheta_2 = \frac{n_1}{n_2} \sin \vartheta_1$$

no solution for
 $n_2 < n_1 \cos \alpha$

total external reflection

critical grazing angle θ_c

$$\theta_c = (4\pi r_o \lambda^2 n)^{1/2}$$

r_o classical e⁻ radius

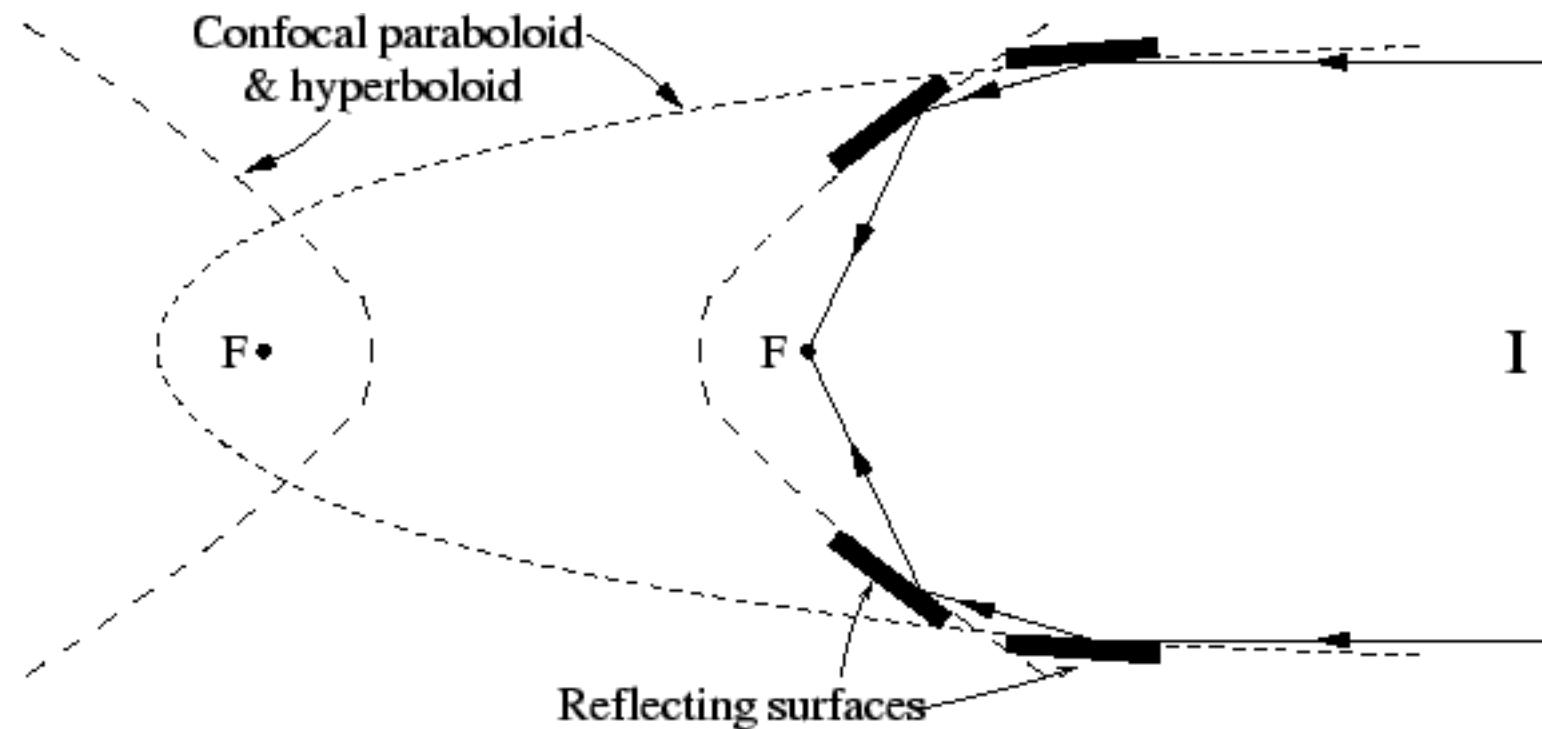
n electron density

- + high reflectivities
- + imaging capabilities
- + large bandpass
- ± small θ_c small (γ -rays) => long focal length
- small projected reflecting surfaces

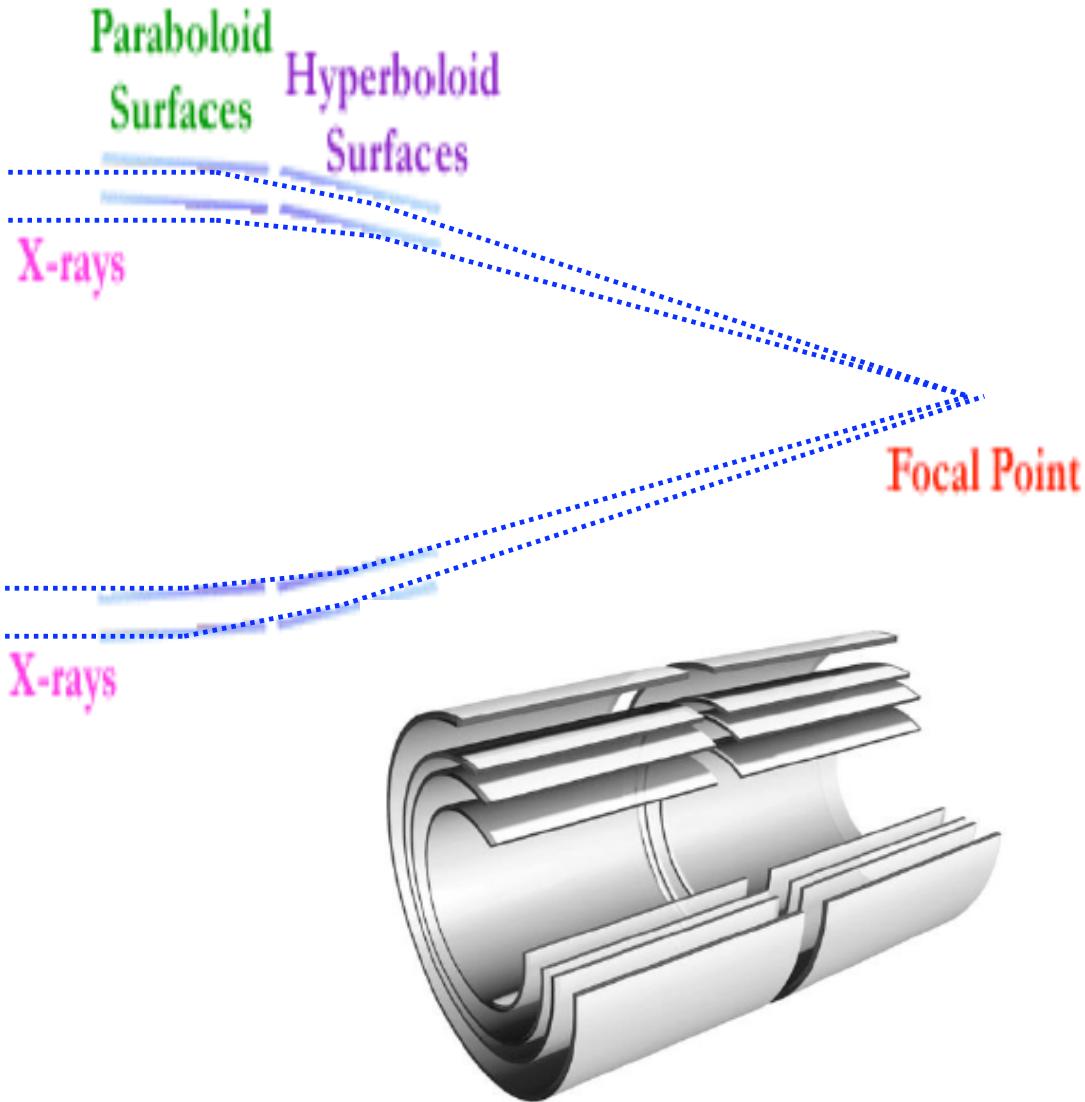
element	Z	δ	β	θ_c (10 keV)
C (diamond)	6	$4.6 \cdot 10^{-6}$	$4.5 \cdot 10^{-9}$	0.173°
Si	14	$4.9 \cdot 10^{-6}$	$7.4 \cdot 10^{-8}$	0.180°
Cu	29	$1.6 \cdot 10^{-5}$	$1.9 \cdot 10^{-6}$	0.326°
Au	79	$3.0 \cdot 10^{-5}$	$2.2 \cdot 10^{-6}$	0.443°

<http://www.astro.caltech.edu/~george/ay20/xray-telescopes.pdf>

Wolter type I telescope

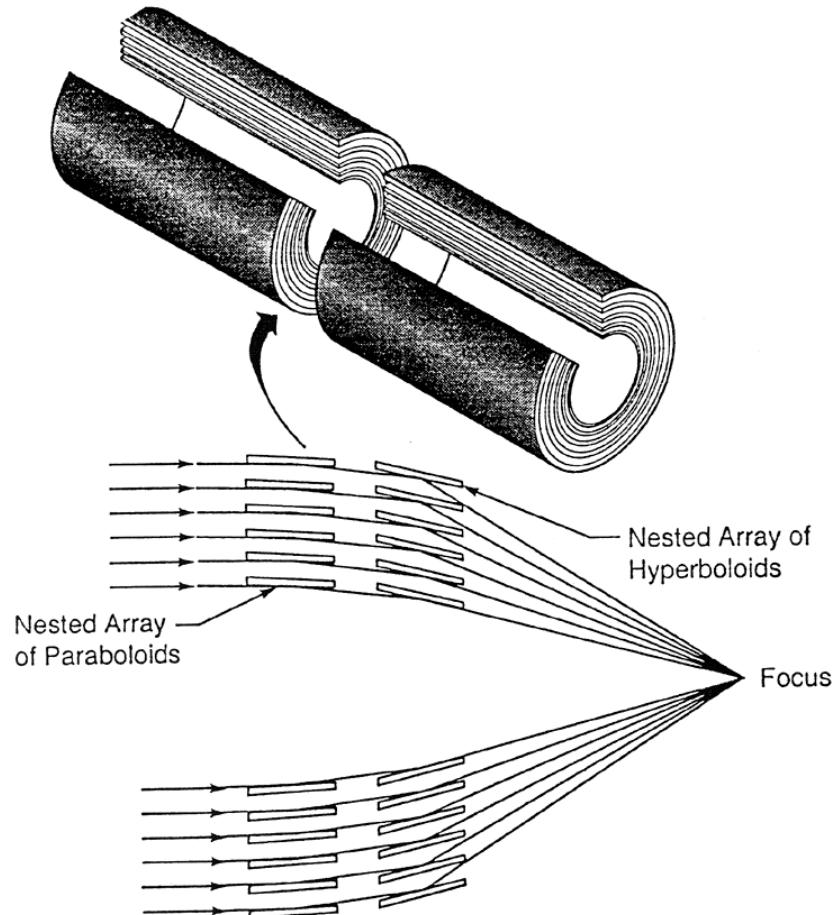


Wolter telescopes



XMM mirror

Wolter telescopes



mission	no	geom	graz.	highest
launch	mirr.	area	angles	energy
		[cm ²]	[arcmin]	[keV]

"traditional" grazing inc. optics (e.g Zerodur)

Einstein '78	4	412	40-70	5
Rosat '90	4	1140	83-135	2
Chandra '99	4	1100	27-51	10

replicated optics

Exosat	2	80	90-110	2
Sax*	4*30	176		
Newton '99	58	6000	18-40	10

foil optics

Asca*	120	4*558	21-45	12
Suzaku/XRT	175	4*873		12

*2 conic sections=approx. Wolter I optic (small θ)

multiple reflection optics (4 or more)

surface imperfection

domain	lambda	$\lambda/10$
Sub-mm	0.5 mm	50 μm
NIR	2 microns	0.2 μm
Visible	0.5 microns	50 nm
X	1 keV	0.1 nm (~ 1 atome)
Gamma	1 MeV	0.1 pm (10-100 x noyau atomique)

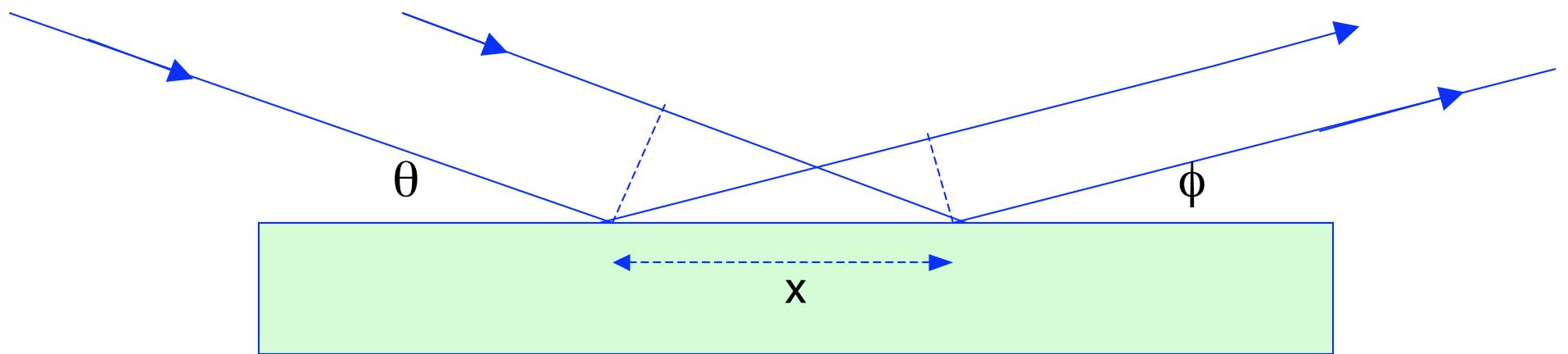
As in an optical telescope, imperfections in the shape of the mirror degrade the resolution.

Good optical surfaces are generally true to about $\lambda/4$ ($\sim 400\text{A}$).

At 10 keV (1 Å), a photon has a wavelength comparable to the size of an atom, and a mirror cannot be polished to levels less than a few atomic diameters.

Fortunately, at grazing incidence, the **effective wavelength seen normal to the mirror is $\lambda/\sin(\theta)$** , so graininess due to individual atoms is not a major concern.

external reflection seen as an interference



$$x \cos \theta - x \cos \phi = n \lambda$$

for all $\lambda \Rightarrow \phi = \theta$

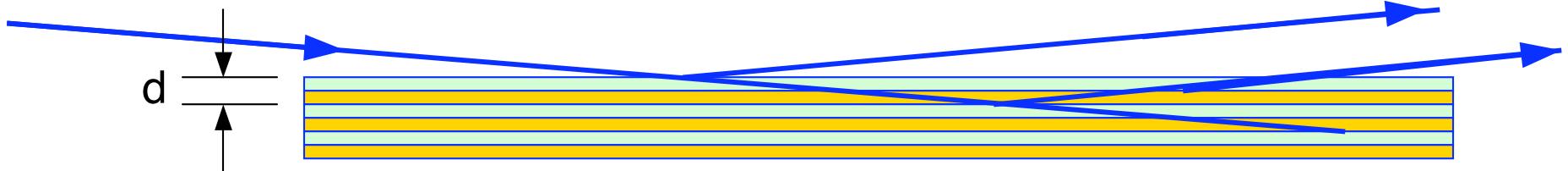
multilayer mirrors at grazing incidence

Wolter-I telescope(paraboloid + hyperboloid)working up to ~ 100 keV

<http://infocus.gsfc.nasa.gov/>

<http://www.srl.caltech.edu/HEFT/index.html>

$$2 d \sin \theta = n \lambda$$



several layers, d constant \Rightarrow narrow bandpass

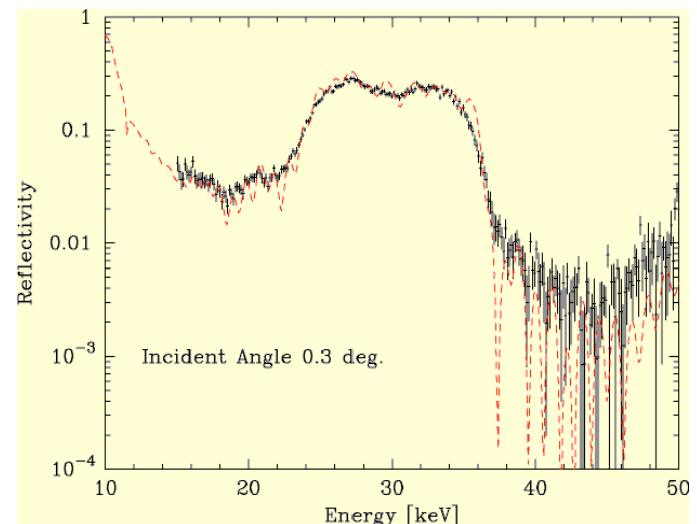
d variable \Rightarrow larger bandpass, reduced efficiency

Example for a multilayer mirror telescope

InFOCuS

Supermirror Hard X-ray Telescope

(GSFC + Nagoya)

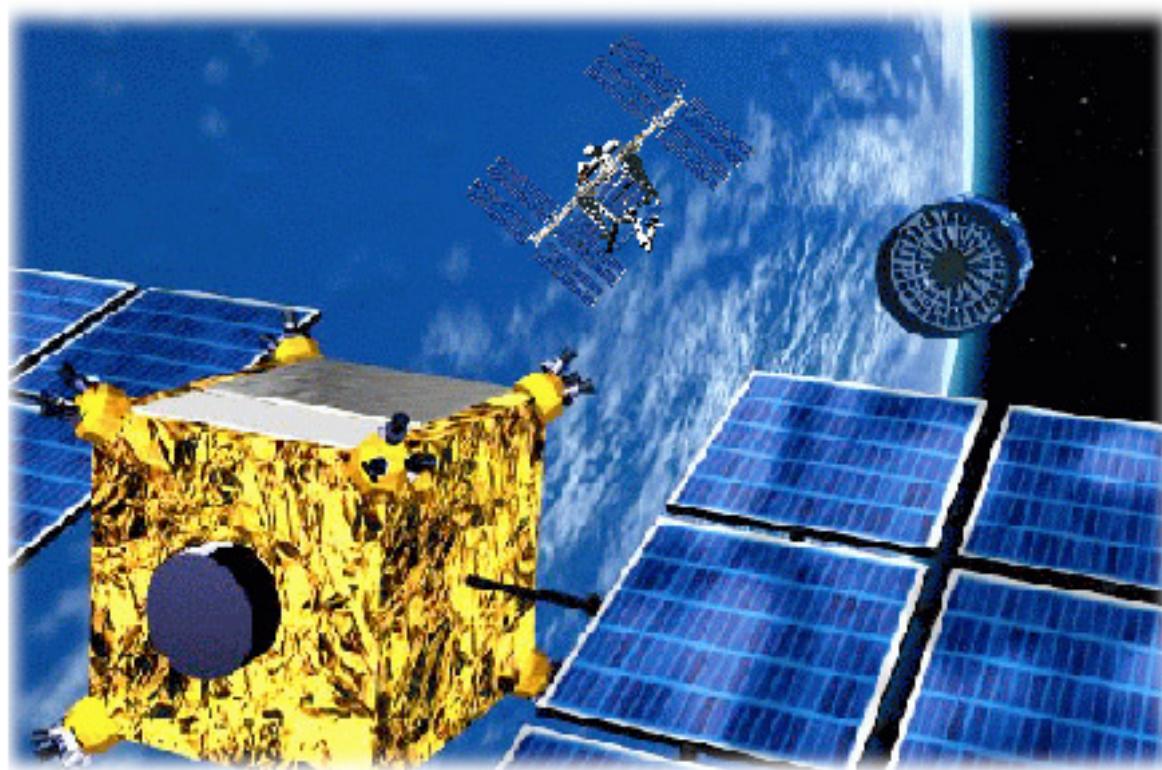


Future : SIMBOLX-X, XEUS

SIMBOLX-X, formation flying mission of CNES (see talk by Ph. Ferrando)

XEUS proposed by/to ESA to succeed XMM-Newton (+ 2015).

- 250 times more sensitive than XMM-Newton
- **30 m²** = (50 x RXTE, XEUS-2)
- 0.05-30 keV (nominal)

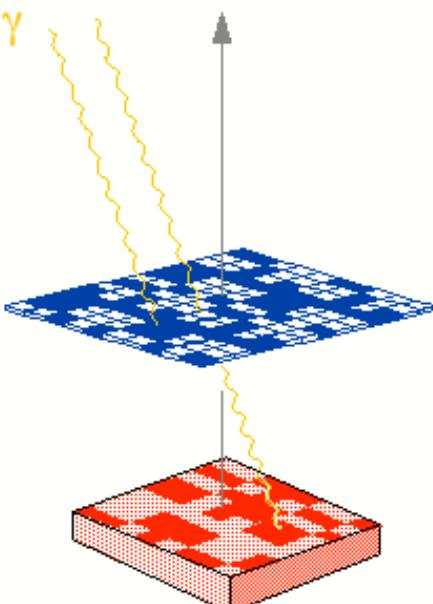
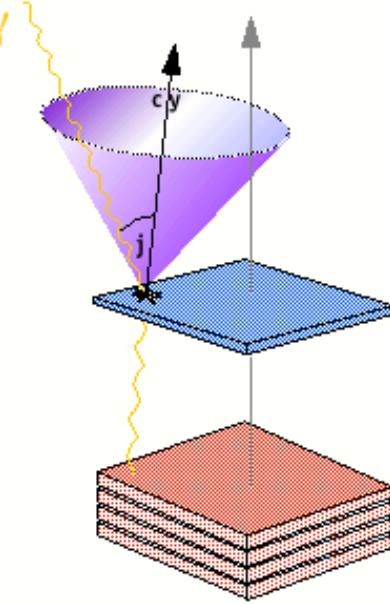
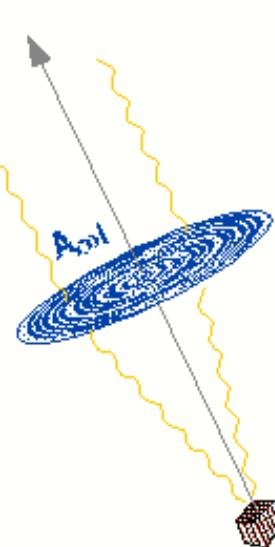


$$N_{\sigma_{\text{qpo}}} \propto \frac{S^2}{(S + B)}$$

S: Source counts
(\propto collecting surface)

B: background counts

Focusing Gamma-Rays - why ?

aperture / effect	modulating aperture systems	Compton telescopes	crystal lens telescopes
	geometric optics absorption	quantum optics incoherent scattering	wave optics coherent scattering
aperture system			
detector	$A_{\text{det}} = A_{\text{col}}$	$A_{\text{det}} = A_{\text{col}}$	A_{det}
signal S	$\sim A_{\text{col}}$	A_{col}	A_{col}
background B	$\sim V_{\text{det}} \sim A_{\text{det}} = A_{\text{col}}$ $\text{const}(A)$	$V_{\text{det}} \sim A_{\text{det}} = A_{\text{col}}$ $\text{const}(A)$	$V_{\text{det}} \sim A_{\text{det}} \ll A_{\text{col}}$ $A_{\text{col}}/A_{\text{det}}$
S/B			

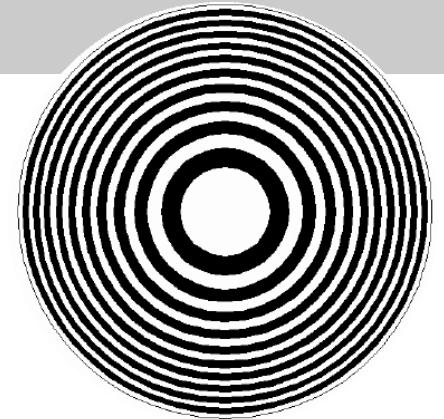
© PvB 1999

Focus Gamma-rays : How ?

I) Refraction / Diffraction in a Fresnel Lens

real part of the refraction index for gamma-rays :

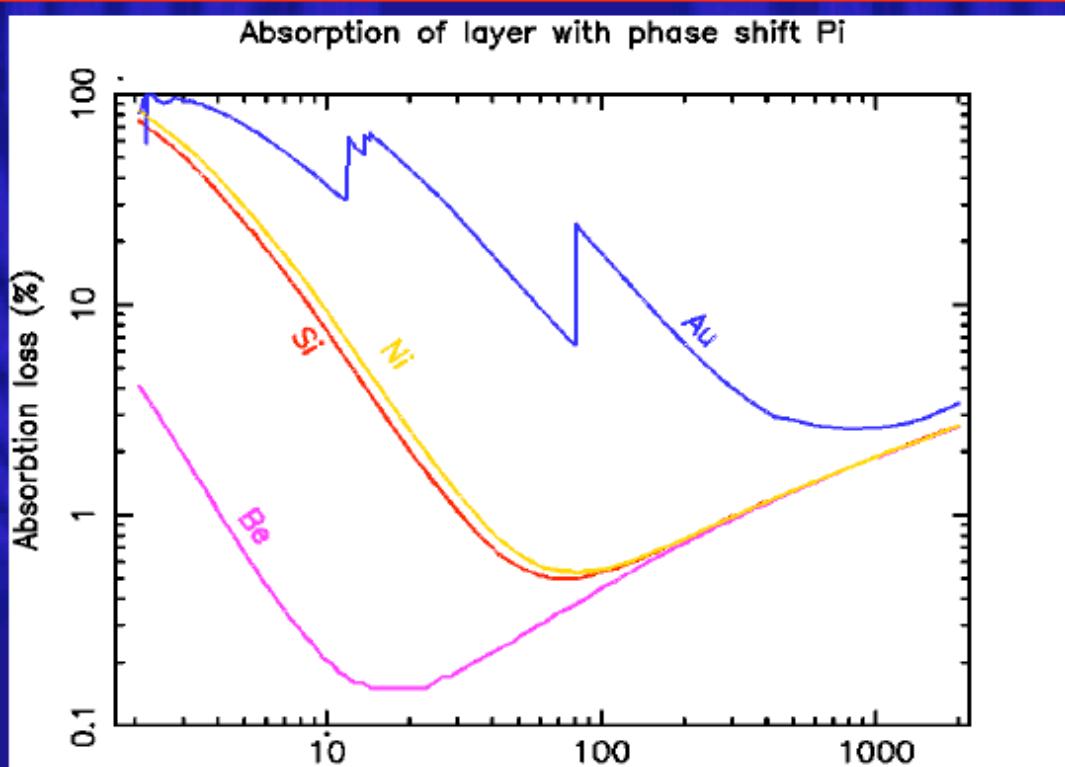
$$\mu = 1 - \delta \quad \text{where} \quad \delta \approx 2 \times 10^{-9} \left(\frac{\rho}{10 \text{ g cm}^{-3}} \right) \left(\frac{E}{1 \text{ MeV}} \right)^{-2}$$



$$t_{2\pi} = \frac{\lambda}{\delta} \approx 0.6 \left(\frac{\rho}{10 \text{ g cm}^{-3}} \right)^{-1} \left(\frac{E}{1 \text{ MeV}} \right) \text{ mm}$$

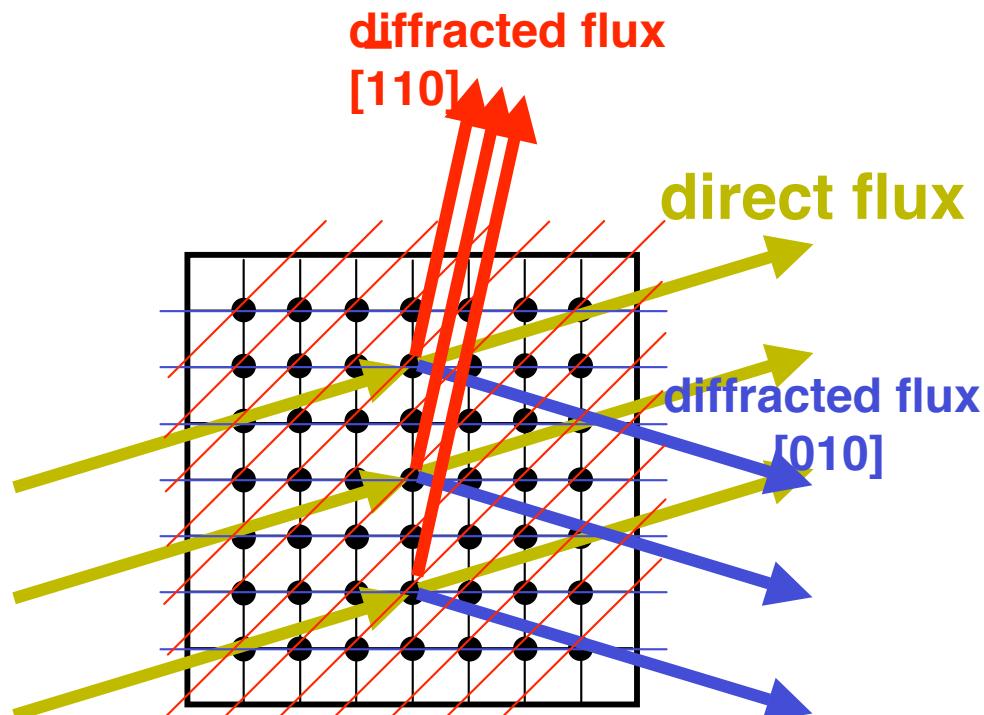
Some examples of the thickness necessary for a phase change of 2π :

	Energy	$t_{2\pi}$	Transmission
Plastic	6 keV	30 microns	96 %
Aluminium	100 keV	225 microns	99 %
Titanium	500 keV	0.7 mm	97 %

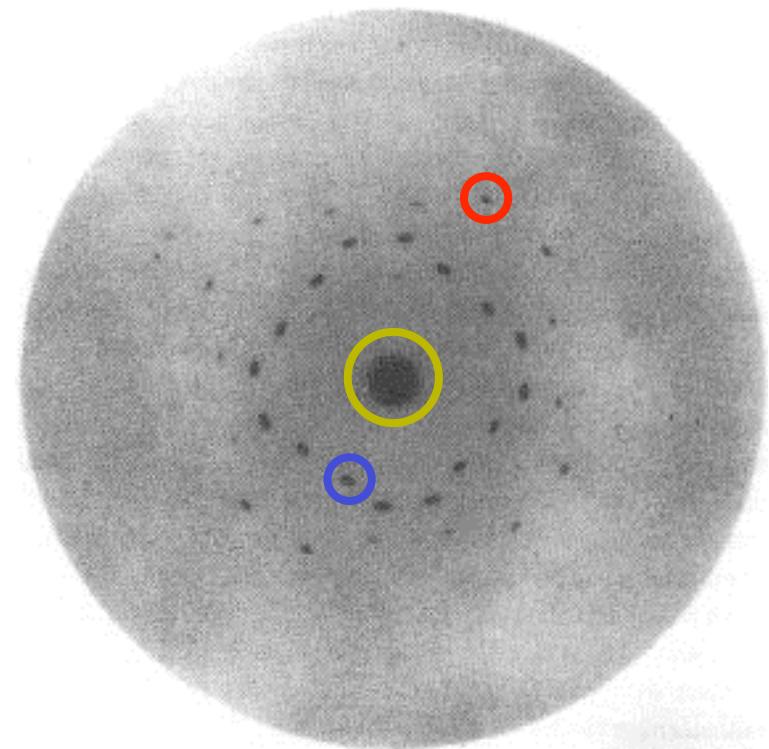


Focus Gamma-rays : How ?

2) Bragg diffraction in a crystal lens

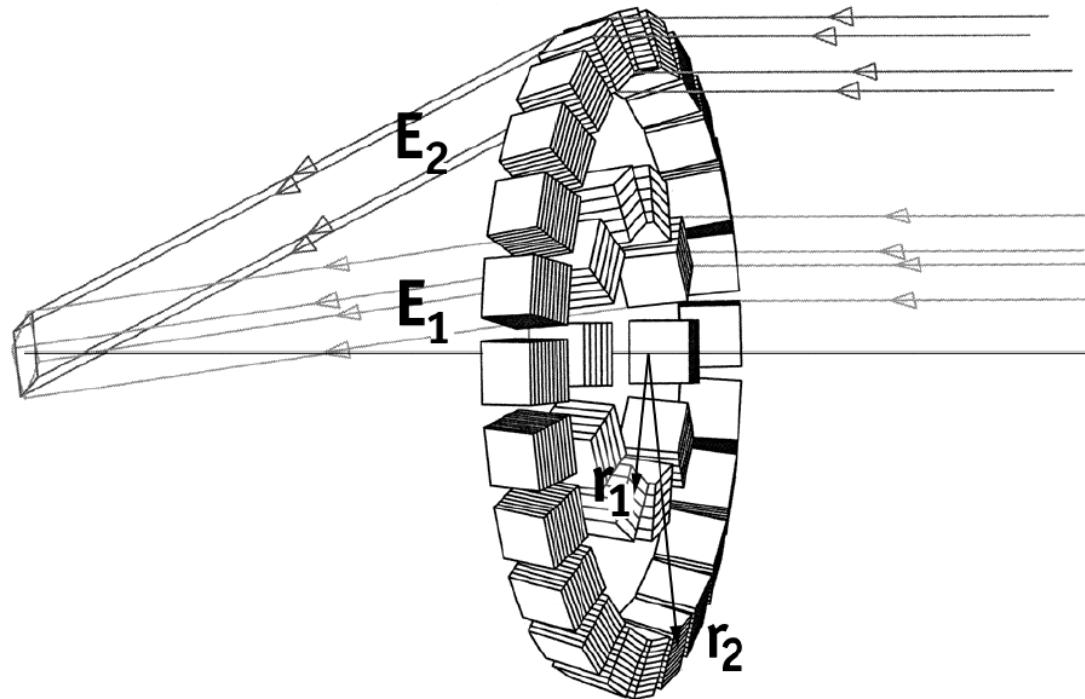


$$2d_{hkl} \sin \theta = n\lambda$$



Laue, Friedrich et Knipping, 1912

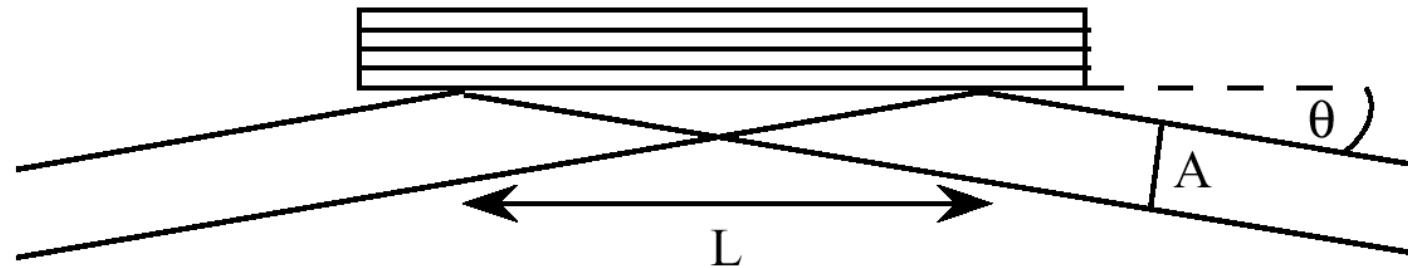
Lens geometry



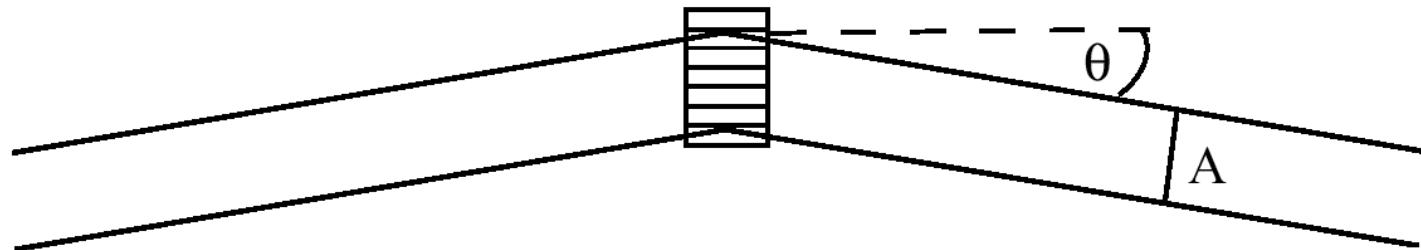
$$2d_{hkl} \sin \theta = n\lambda$$

Bragg vs. Laue geometry

Bragg geometry



Laue geometry



Laue geometry uses volume reflection, whereas Bragg geometry uses surface reflection.

Bragg geometry : a beam with cross-section A would require a crystal length L
 $L = A / \sin\Theta_B$ e.g. for $A=1\text{cm}$ and $\Theta_B=0.5^\circ$ \Rightarrow $L = 114\text{ cm}$

calculating the reflection efficiency of a Laue lens

simplified calculation using the Darwin mosaic model :

- reflection amplitudes of each microcrystal added incoherently
- absorption is calculated on a macroscopic scale
(the absorption within each microcrystals is neglected)

The reflection efficiency at the Bragg energy as a function of Bragg angle θ is:

$$r_{th}(\theta) = 0.5(1 - e^{-2\alpha T})(e^{-\mu T})$$

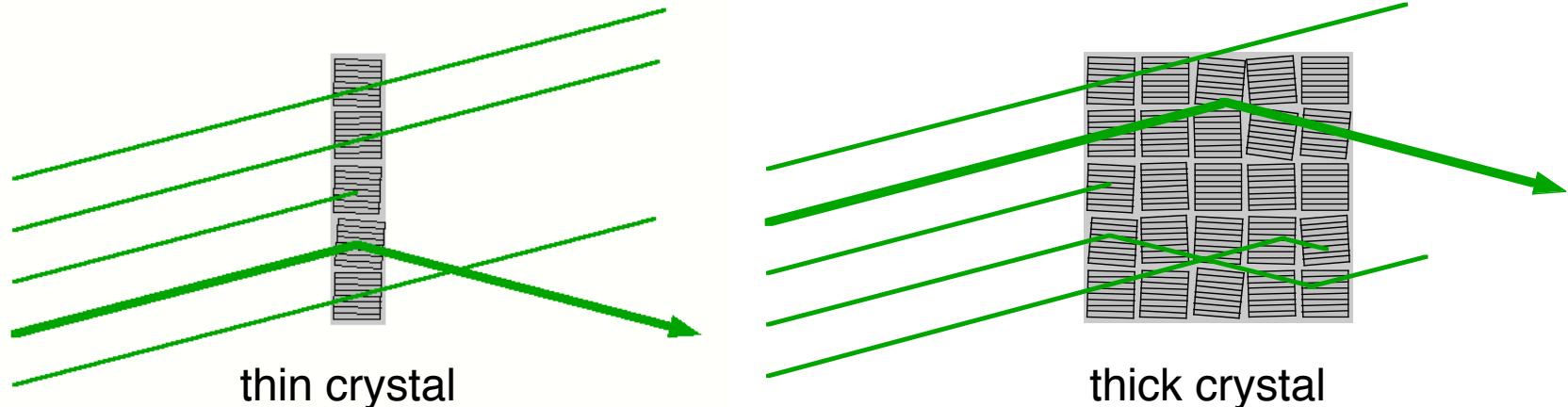
μ Absorption coefficient (T crystal thickness)

$\alpha(\theta)$ diffraction coefficient : $\alpha(\theta) \sim \frac{F^2 \lambda^3}{V^2 \sin(2\theta)} \sim \frac{\rho^{5/3}}{E_\gamma^2}$

F : structure factor $F[hkl=333] \ll F[hkl=111]$

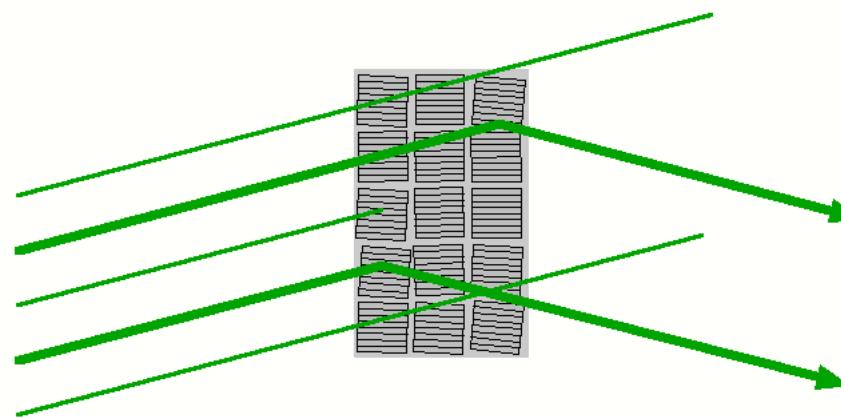
The reflection efficiency decreases with increasing energy and order,
and with decreasing structure factor.

the optimal thickness for a diffraction crystal (E)



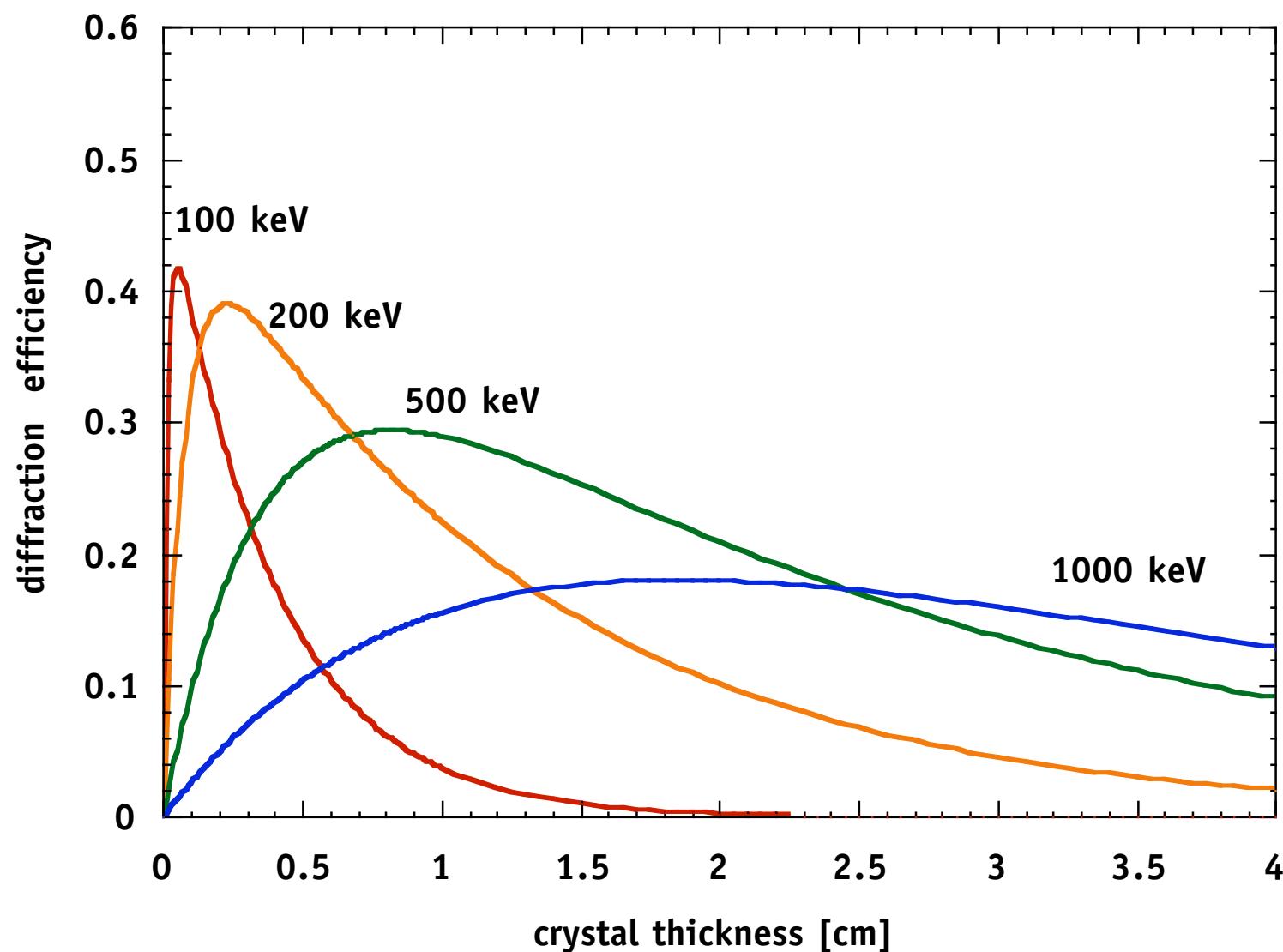
thin crystal

thick crystal

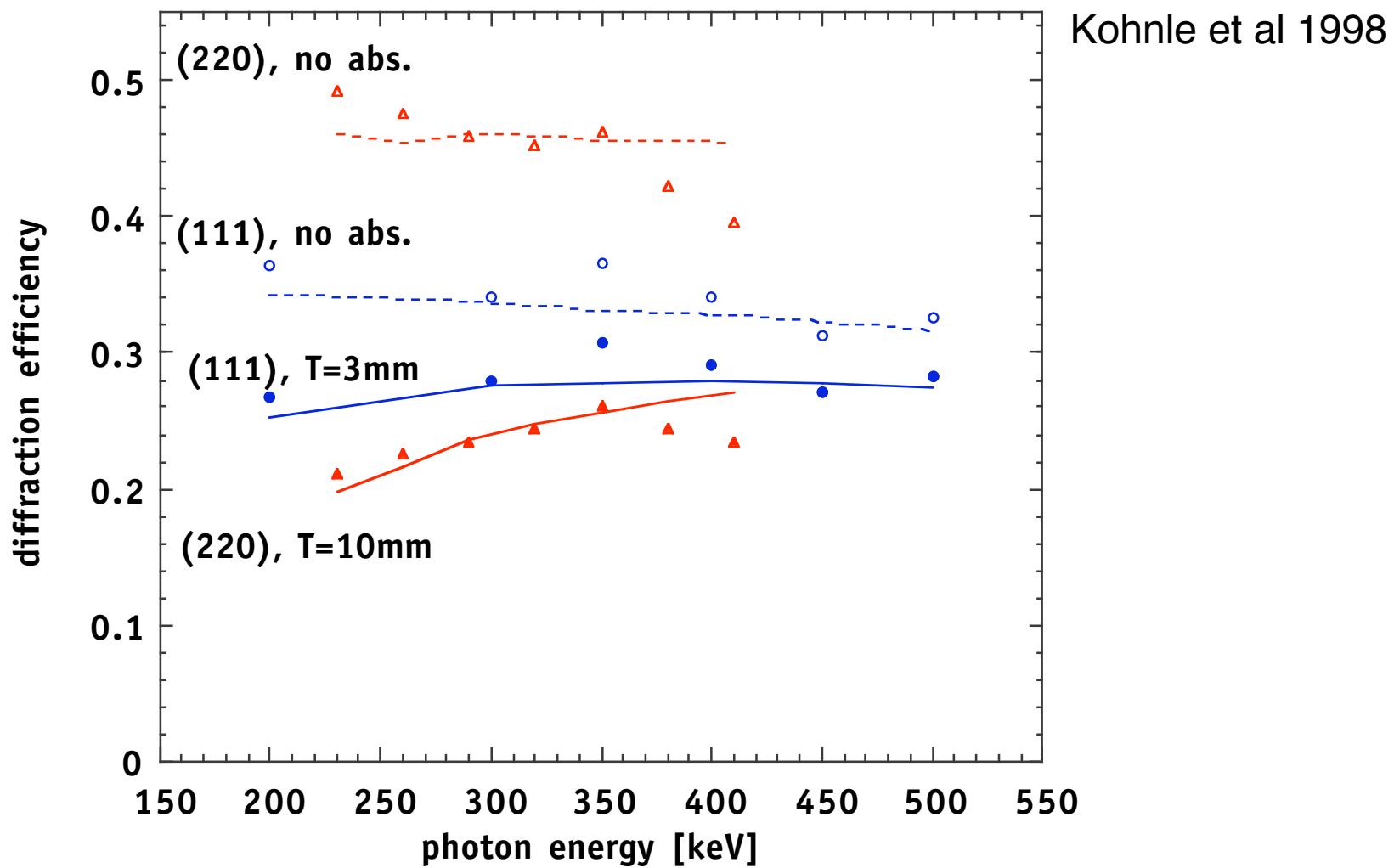


optimal thickness - energy dependent !

the crystal diffraction efficiency for Ge [440] planes

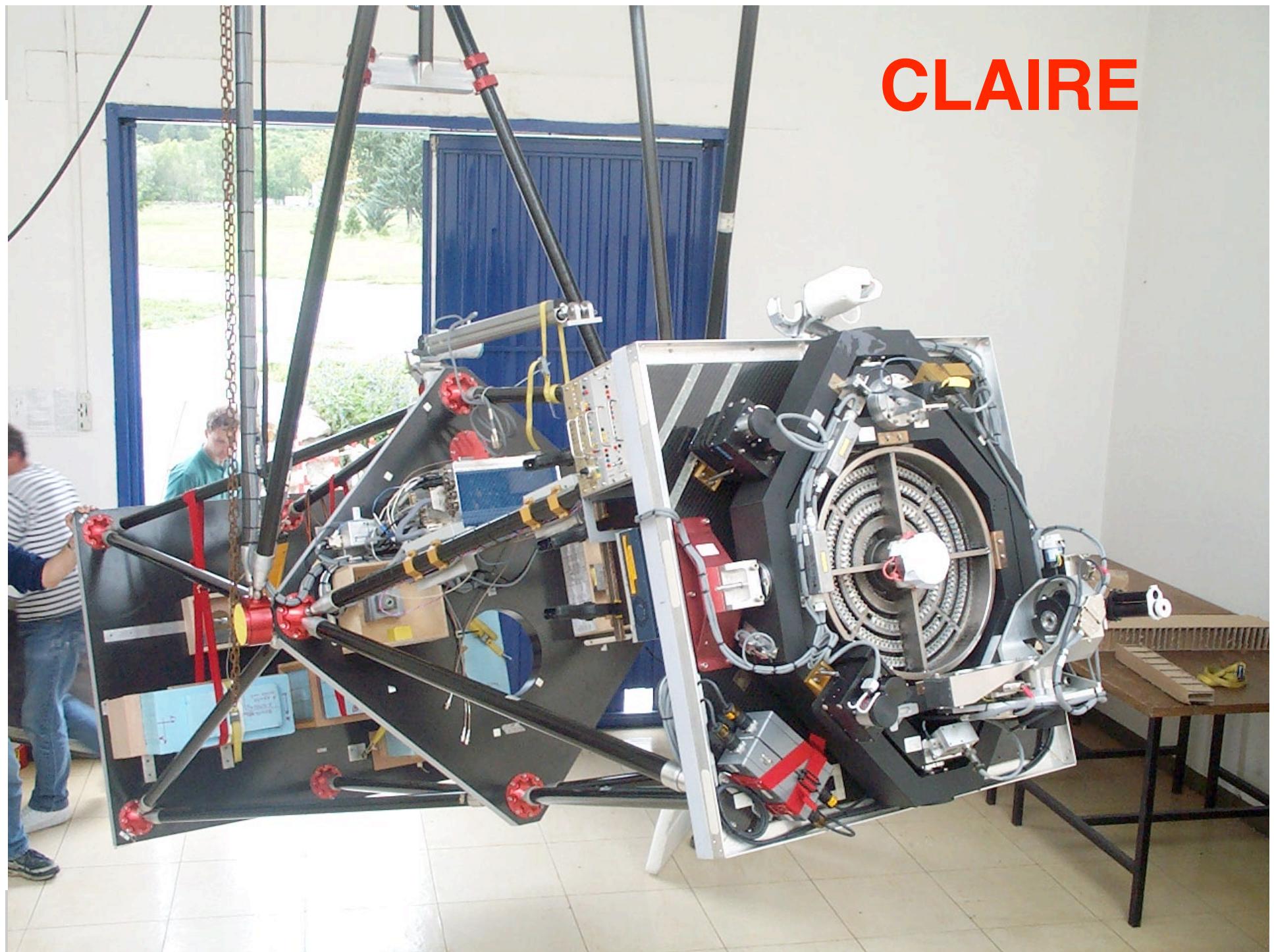


Diffraction efficiency of Ge crystals

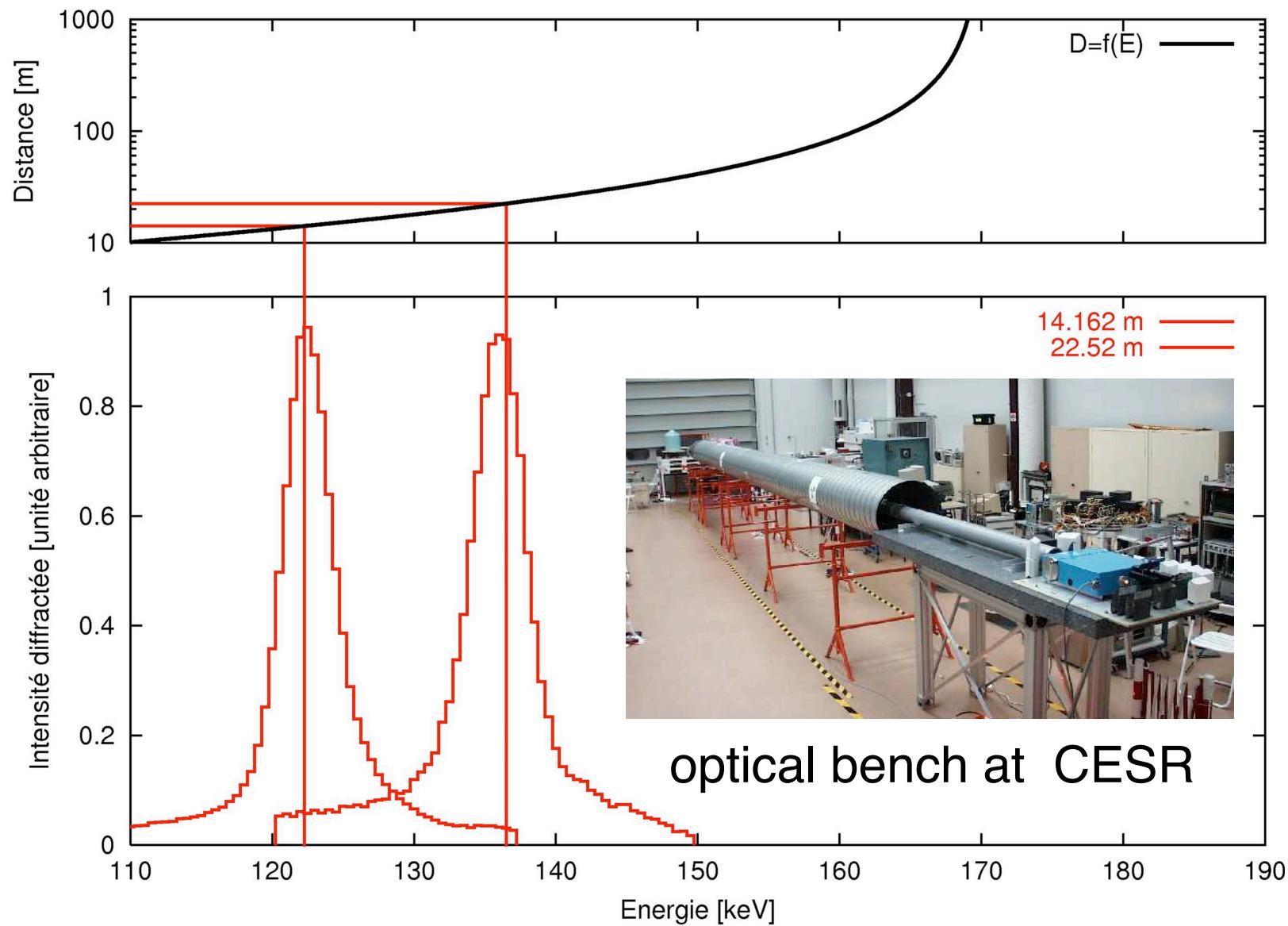


APS / ANL synchrotron beam (\sim parallel, divergence $\approx 3''$)
diffraction efficiency : ratio of doubly diffracted / singly diffracted flux

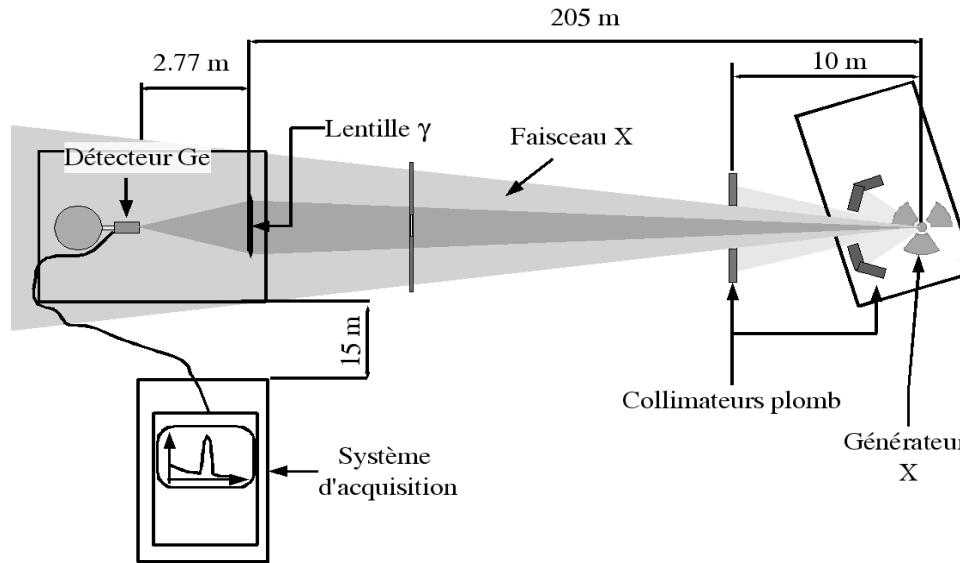
CLAIRES



CLAIRE : tests in the lab ... and beyond

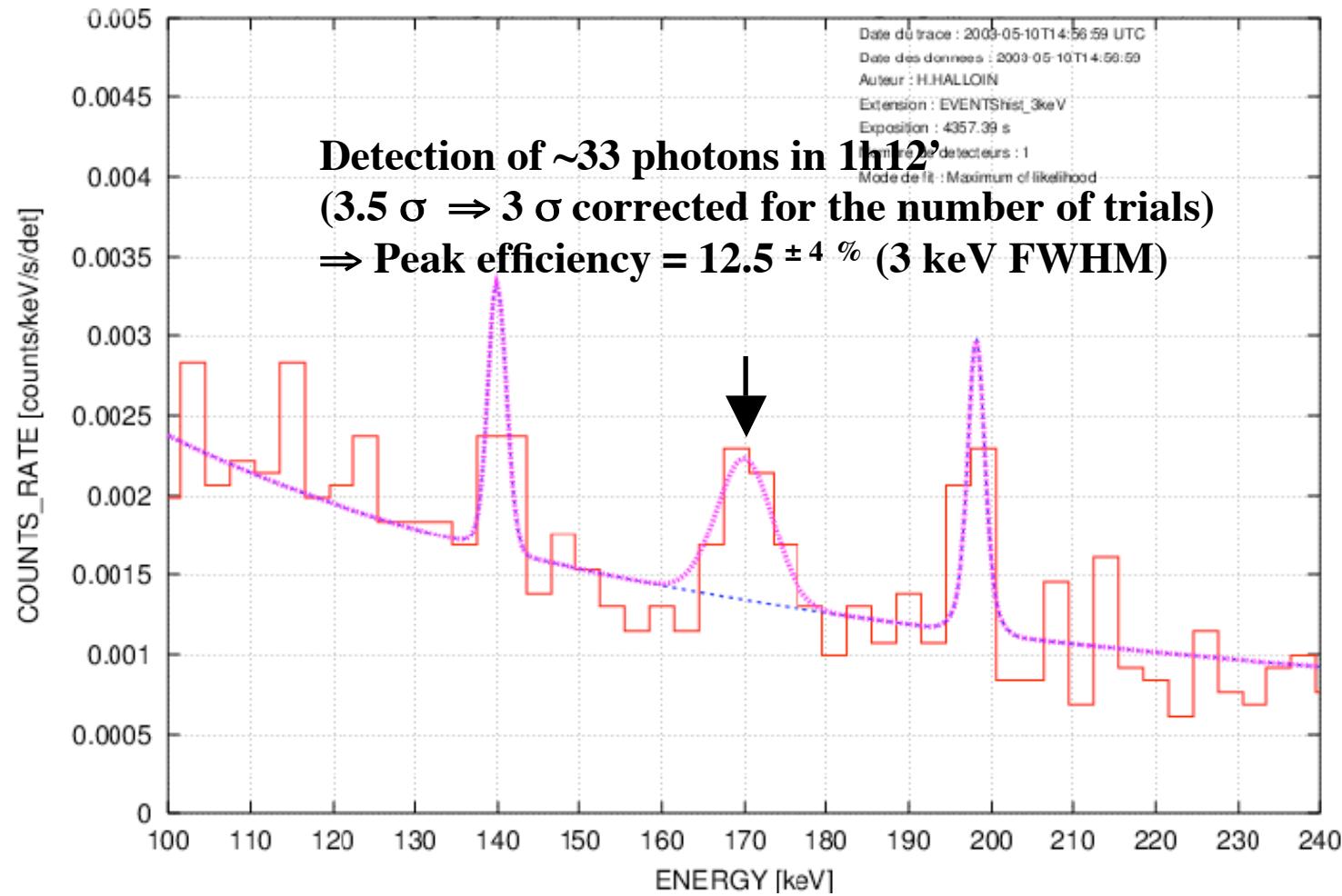


CLAIRE TGD : a source close to infinity ...

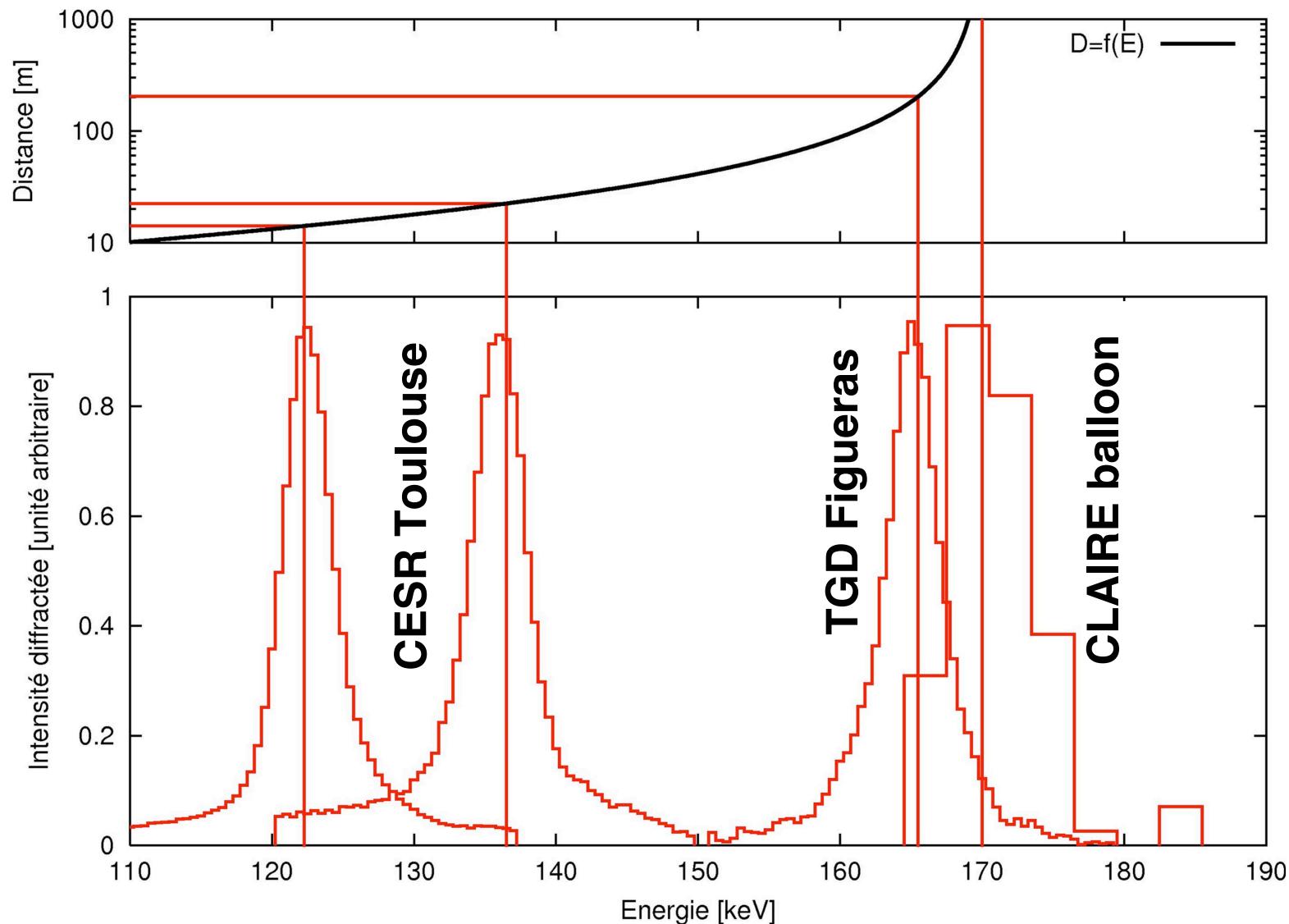




CLAIRE 2001 : first light for an astrophysical source

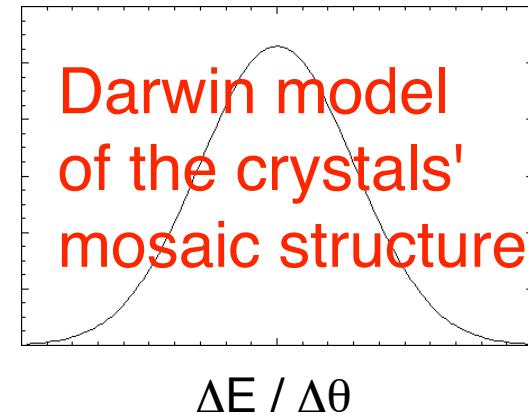
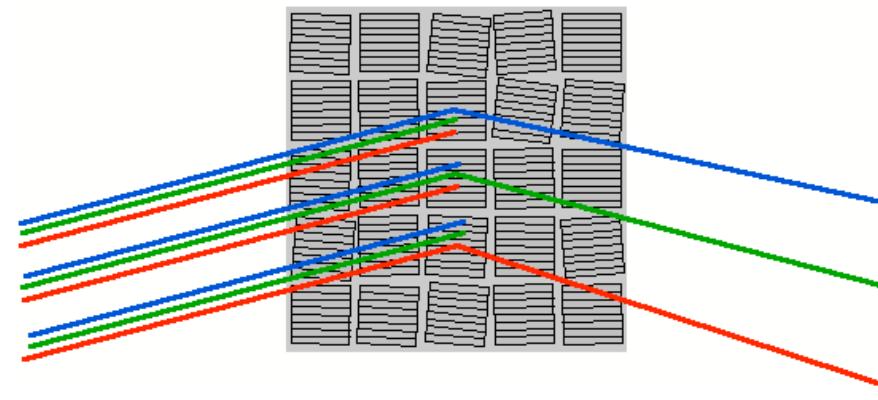
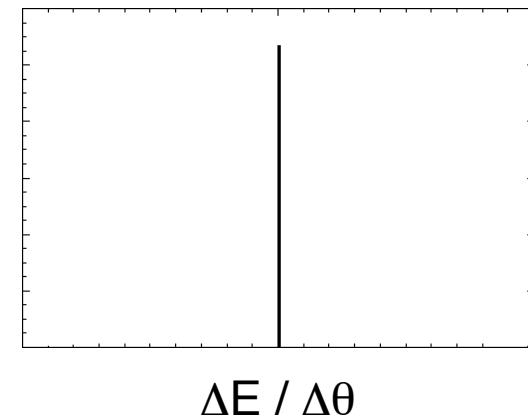
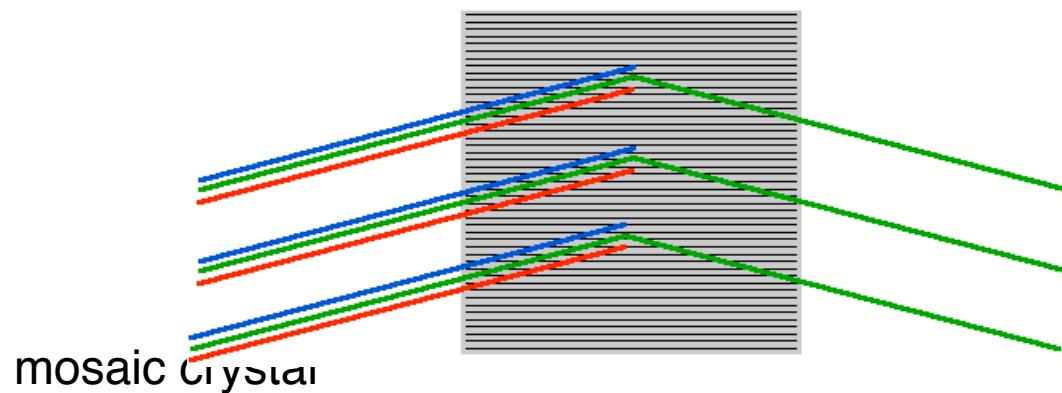


CLAIRE : 14 m, 22.5 m, 205 m ... infinity ! $\epsilon_{\text{peak}, 3 \text{ keV}} \approx 10 \text{ \%}$

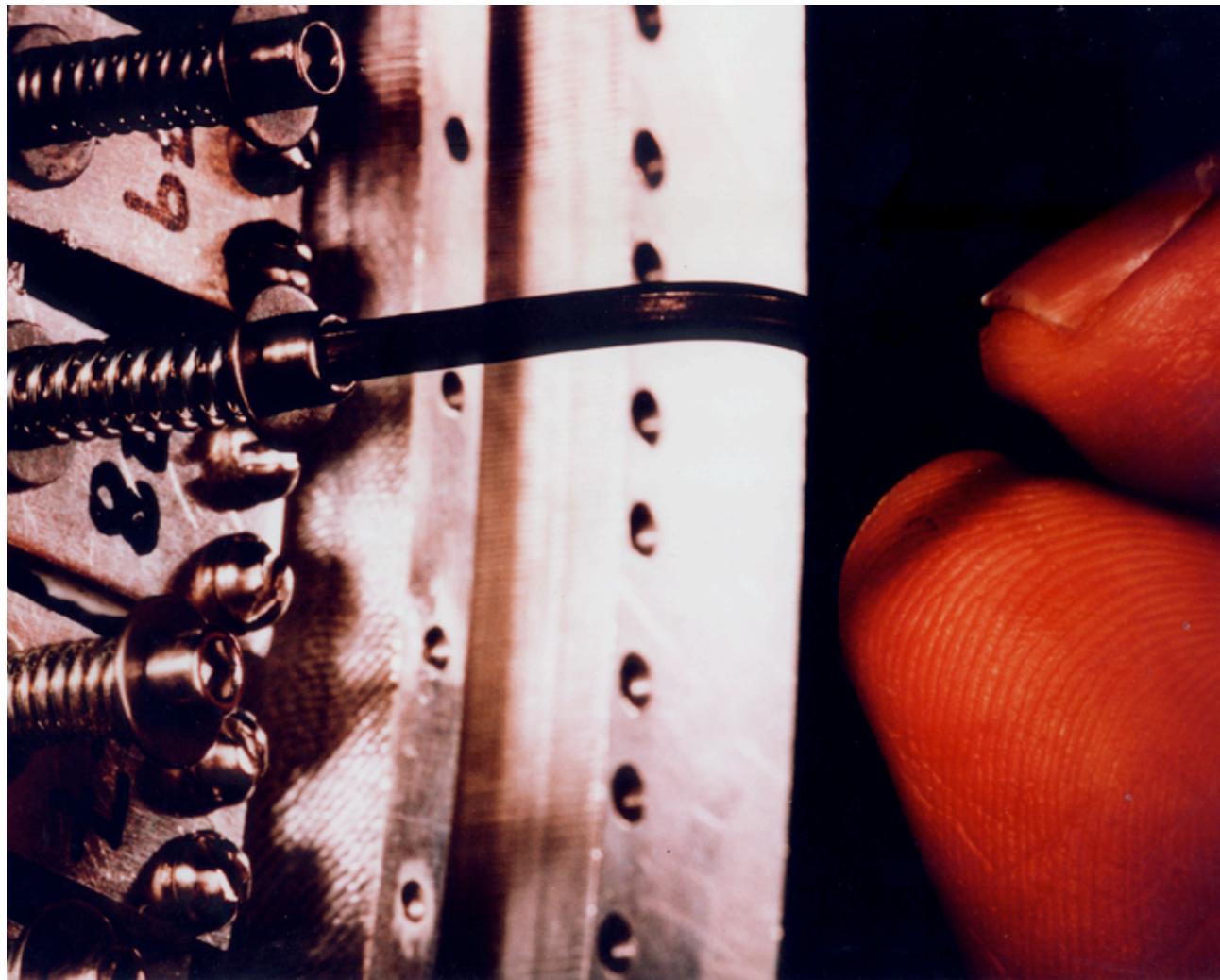


the energy bandpass ΔE and field of view $\Delta\theta$

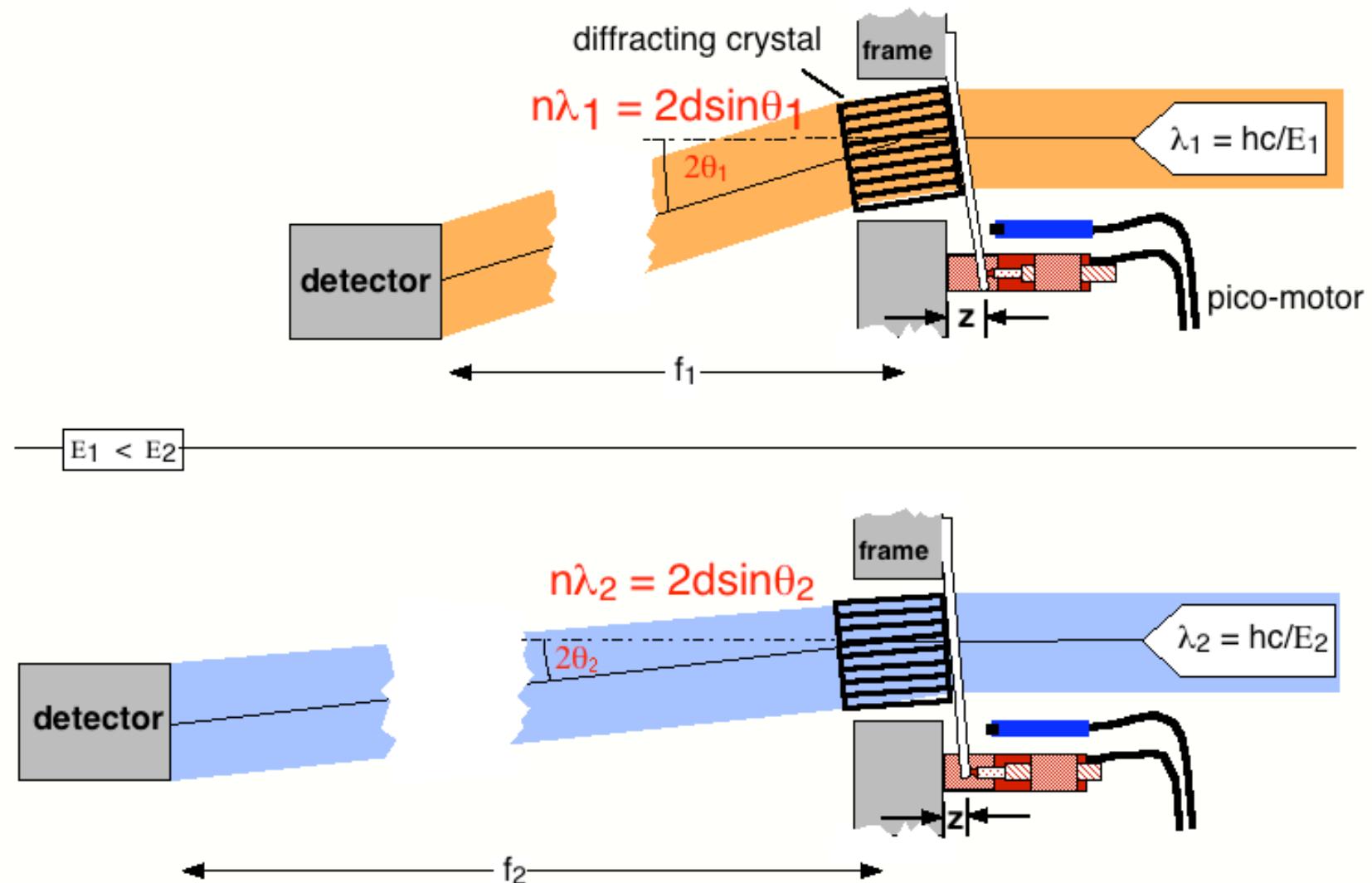
perfect monocrystal



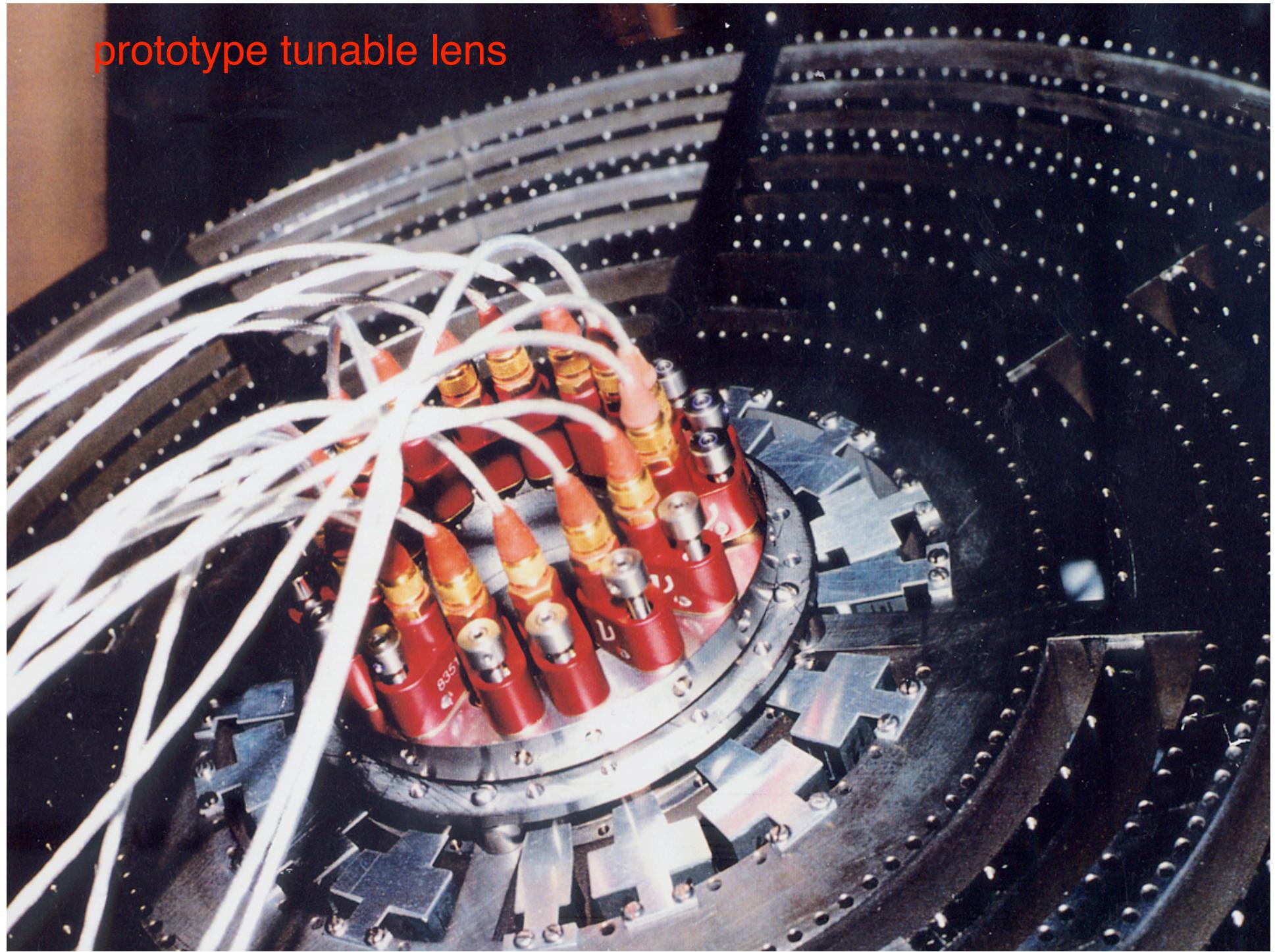
tuning the lens ...



the principle of a tunable Laue Lens



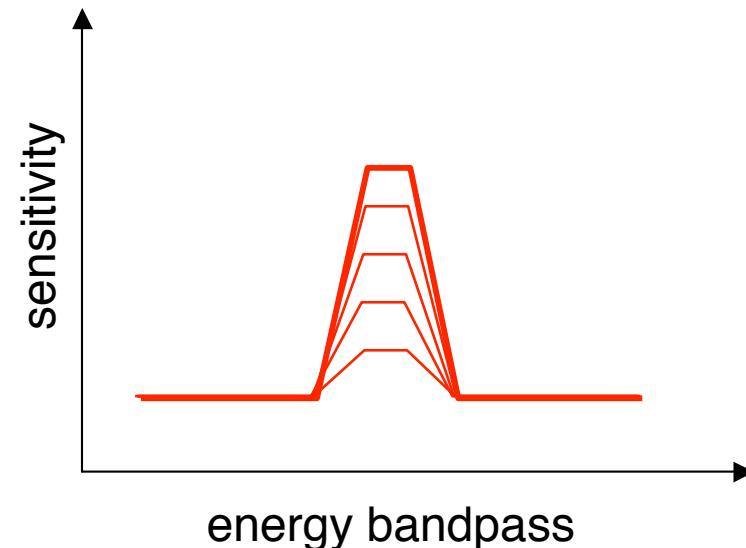
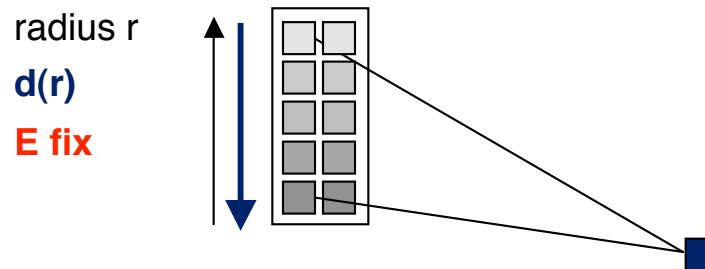
prototype tunable lens



everything about nothing or nothing about everything

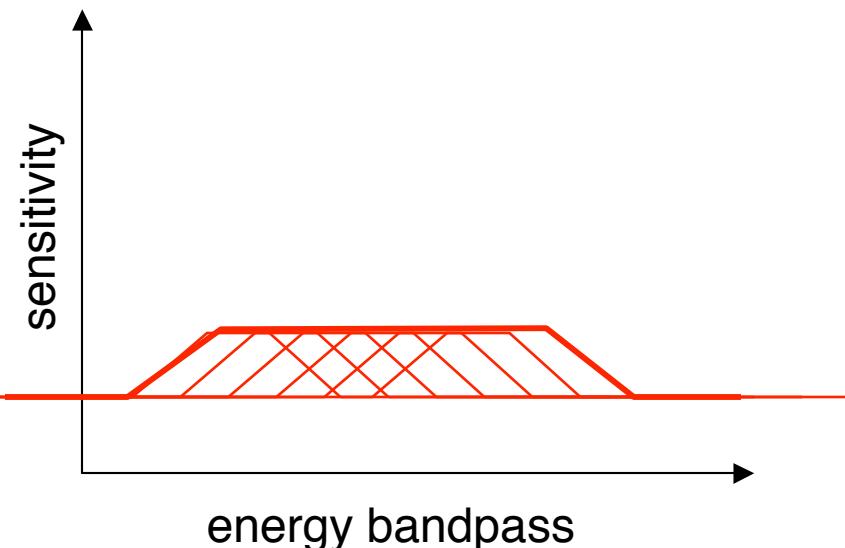
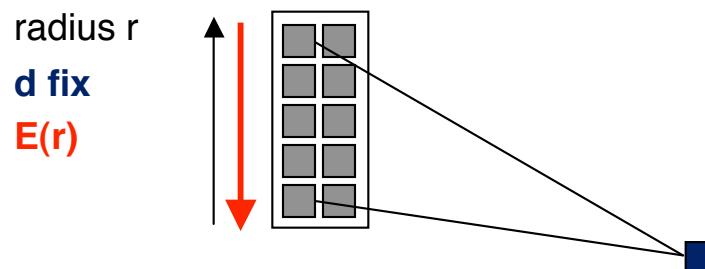
Narrow band Laue Lens :

Each ring uses **different set of crystal planes**

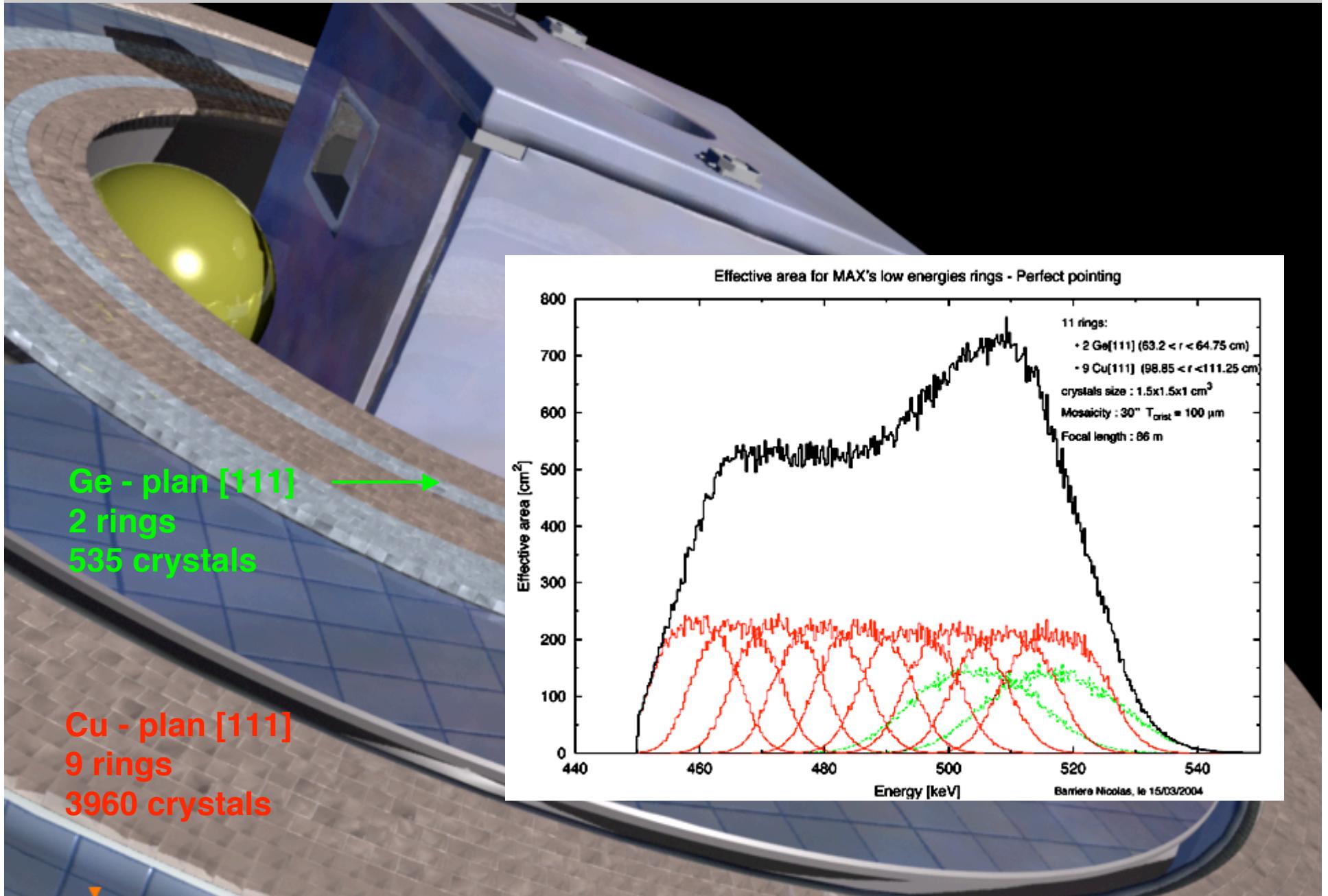


Broad band Laue Lens :

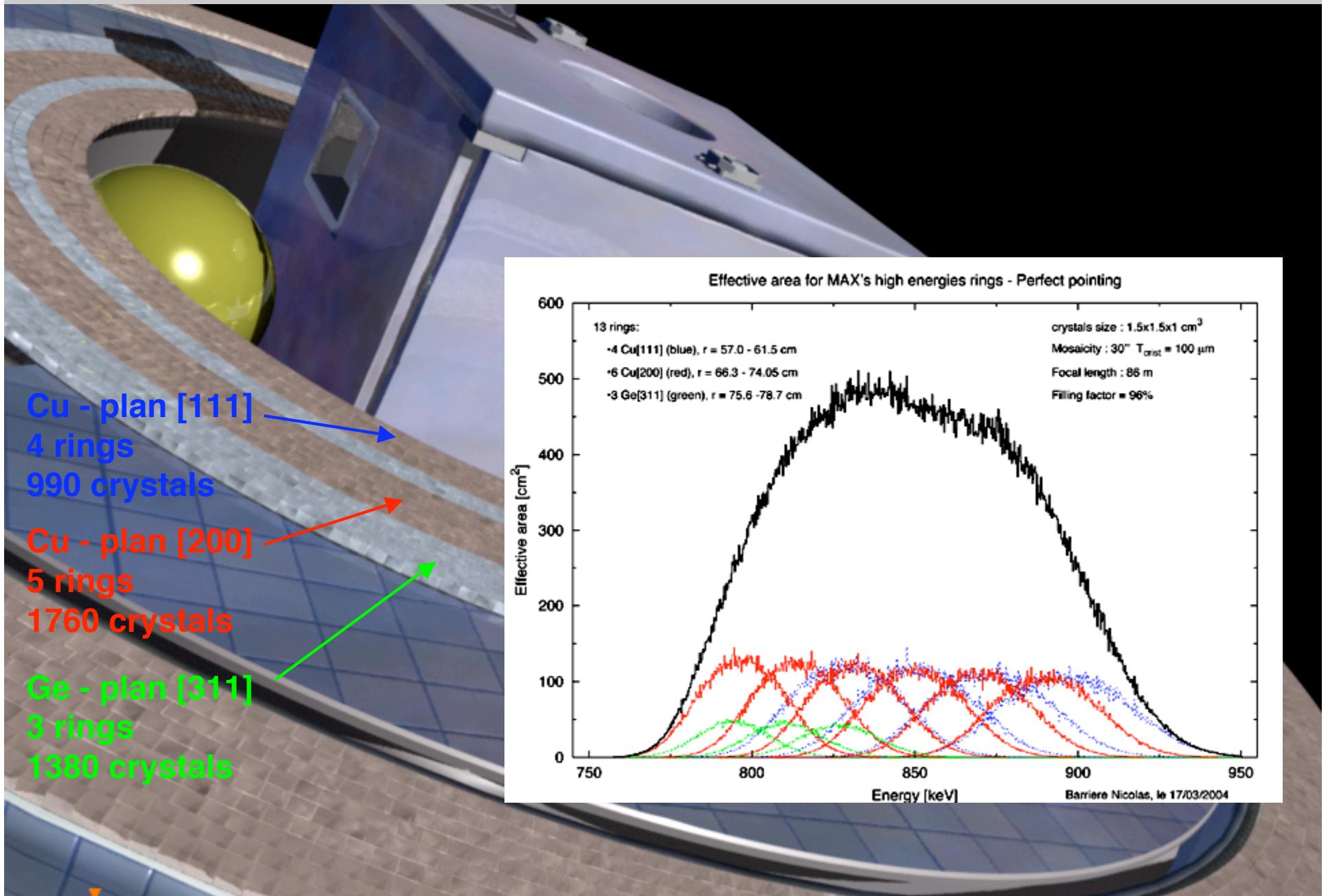
All rings use **same set of crystal planes**

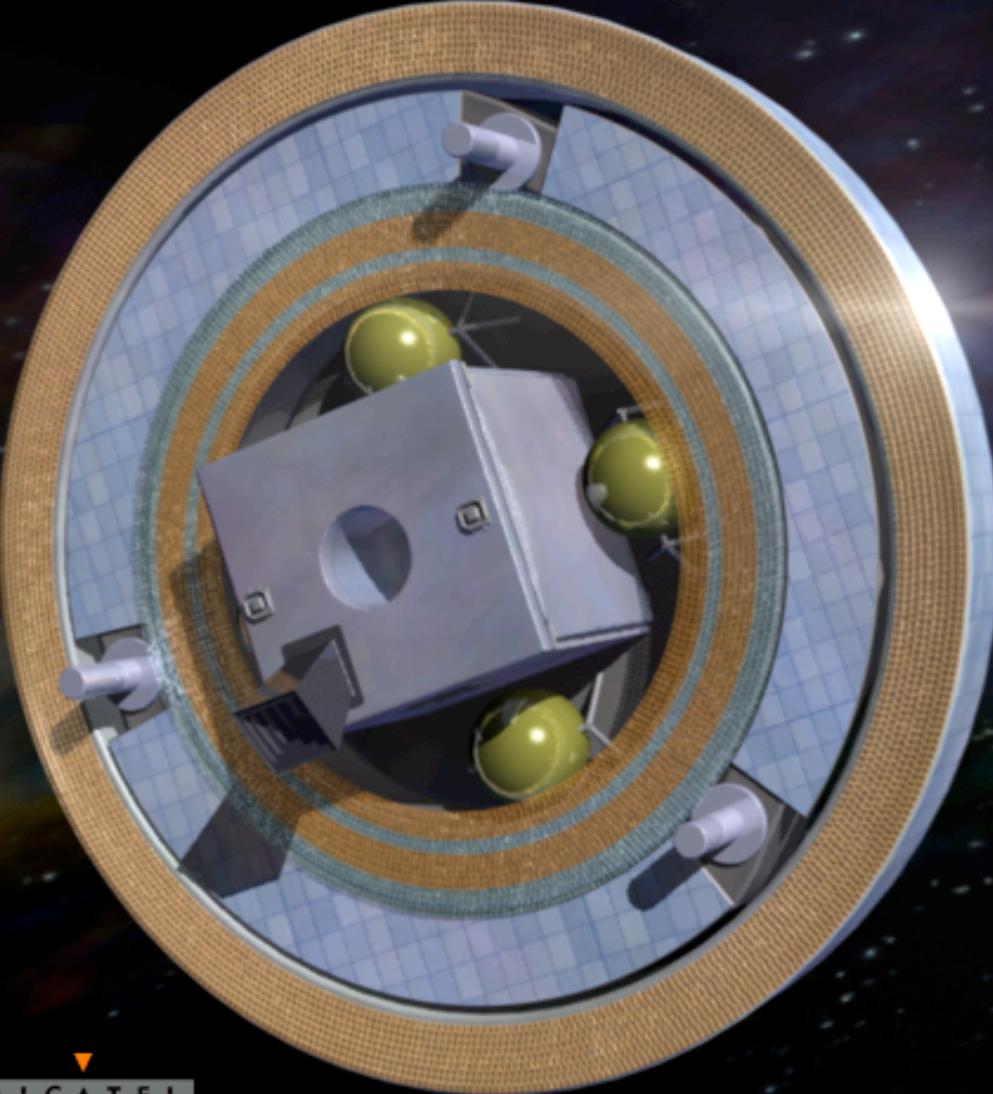


designing a Laue lens telescope - efficient area I



designing a Laue lens telescope - efficient area I



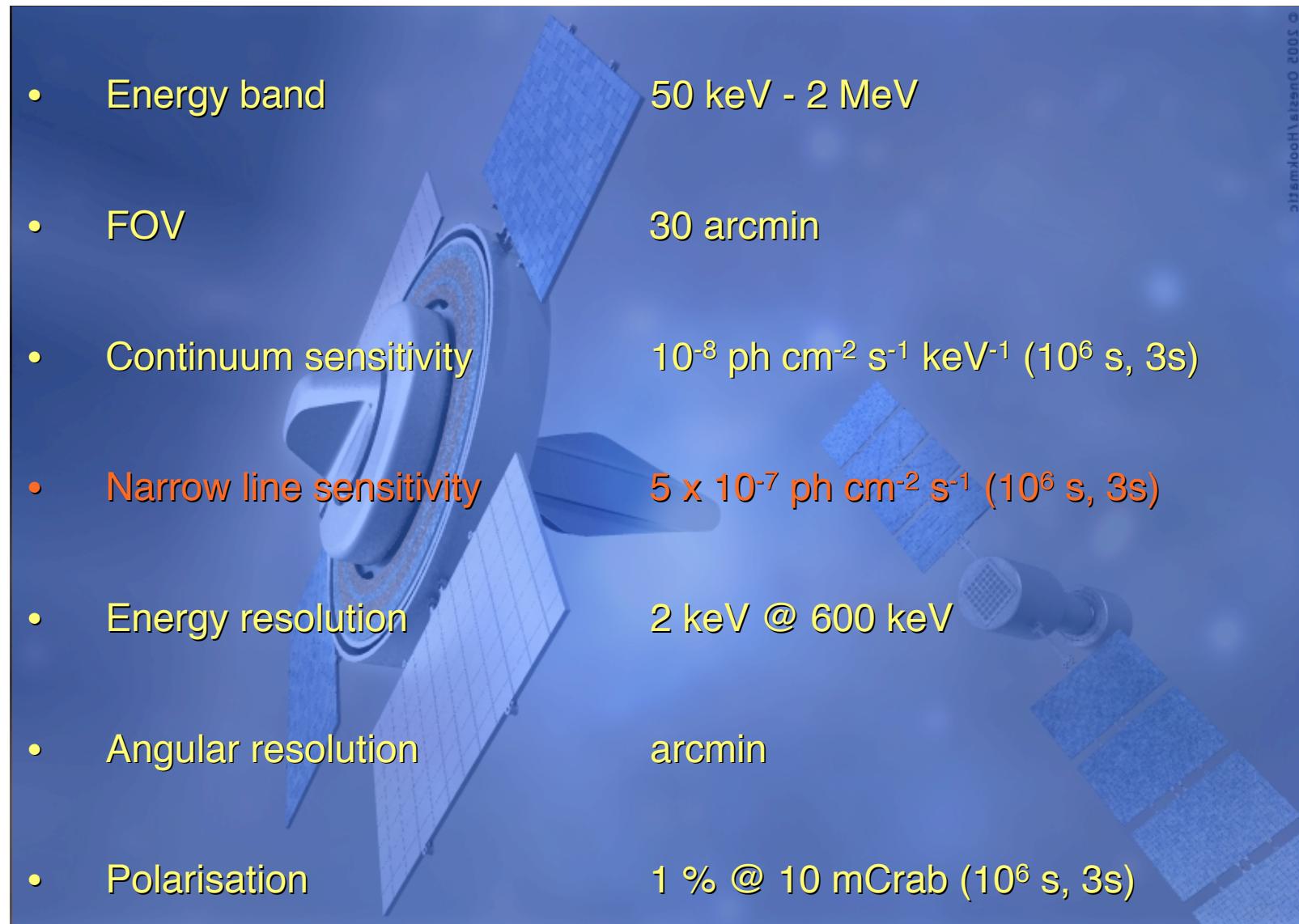


MAX



ALCATEL
SPACE
David Biau 2004

Mission requirements for a Gamma-Ray Imager



• Energy band	50 keV - 2 MeV
• FOV	30 arcmin
• Continuum sensitivity	$10^{-8} \text{ ph cm}^{-2} \text{ s}^{-1} \text{ keV}^{-1}$ ($10^6 \text{ s}, 3\text{s}$)
• Narrow line sensitivity	$5 \times 10^{-7} \text{ ph cm}^{-2} \text{ s}^{-1}$ ($10^6 \text{ s}, 3\text{s}$)
• Energy resolution	2 keV @ 600 keV
• Angular resolution	arcmin
• Polarisation	1 % @ 10 mCrab ($10^6 \text{ s}, 3\text{s}$)